TABLE 2. Values of the Interplanar Distances of Zirconium under a Pressure P = 350 kbar, and the Interplanar Distances of Its Modifications

Sample under pres- sure P = 350 kbar		Original structure [7]			Structure of the w phase [4]			Structure of new phase		
d	I	d	hkl	I	d	hkl	I	d	hkl	I
2.794 2.568 2.523 2.460 1.897 1.784 1.615 1.395 1.366 1.350 	med med med w.s. s	2.798 2.573 — 2.459 1.894 1.616 1.463 1.399 1.368 1.350 — 1.229 1.0842 1.0360	010 002 011 110 110 013 020 112 021 1- 004 022 023 121	33 33 33 	3.103 — 2.533 2.521 1.957 1.786 1.555 1.455 — 1.323 1.317 1.259 1.129 1.062 1.037 1.029 1.008	001	6 — 45 100 13	1.783 	011 	1000
0.9783 0.9534 —	w. s. —	0.9783	123	2 _ _	0.9784 0.9525 0.9517	222 231 140	15 13 15	0.9534	123	60

Note. We use the following notation: d are the interplanar distances; hkl are the Miller indices; I is the intensity; v.s. is very strong; s is strong; med. is medium; w. is weak; v.w. is very weak.

As the table shows the new high-pressure phase we obtained in zirconium after a shock wave has passed through it is not the ω phase of [4]. This is indicated by the absence on x-ray diffraction photographs of the ten comparatively strong lines of the ω phase, and also the lack of correspondence of experimental and calculated intensities. The indexing of seven lines of the new phase showed that all of them correspond with great precision to a lattice with the structure of a body-centered cube with $\alpha=3.568\pm0.005$ Å. The density of the new phase $\rho=6.656$ g/cm³.

A high-pressure phase was also identified in titanium, but only in samples subjected to a pressure $p=350~\mathrm{kbar}$. On the x-ray diffraction photographs of these samples three new, not very strong lines are present. The absence of lines of a new phase in titanium samples subjected to a pressure $P=500~\mathrm{kbar}$, and their low-intensity in zirconium samples for the same pressures are apparently explained by their instability at the high residual temperatures of shock compression. Just as in zirconium, these lines can be ascribed to the strongest reflections from (011), (002), and (022) of a cubic phase with $a=3.27~\mathrm{Å}$. Unfortunately, these lines coincide with (110), (021), and (220) of the ti-

tanium ω phase. The absence of other strong reflections of the ω phase compels us to presume that we succeeded in isolating the body-centered cubic high-pressure phase in titanium, as in zirconium.

The last column of Table 1 shows the parameters of the new metastable phases of Zr and Ti; they were, however, obtained at much lower temperatures.

LITERATURE CITED

- 1. P. W. Bridgman, Proc. Am. Acad. Arts Sci., 76, 71 (1948).
- 2. P. W. Bridgman, Proc. Am. Acad. Arts Sci., 81, 165 (1952).
- 3. P. W. Bridgman, Proc. Am. Acad. Arts Sci., 76, 55 (1948).
- 4. J. C. Jamieson, Science, 3 (3562), 72 (1963).
- A. Gayaraman, W. Klement, and G. G. Kennedy, Phys. Rev., 131, 644 (1963).
- V. L. Al'tshuler, I. M. Barkalov, I. N. Dulin,
 V. N. Zubarev, T. N. Ignatovich, and P. A.
 Yampol'skii, Khimiya Vysokikh Energii, 2, 88 (1968).
- 7. L. I. Mirkin, Handbook of X-ray Structural Analysis [in Russian], GIFML, Moscow (1961).

ns sam-) shortmplimethod (at of tanium vestiga-

d in zirplitudes. ich exobtained e result zirar. al zirdisphase l also s for a ie table re cal-Lorents. fac-